

## Synthesis, structural determination and magnetic behaviour of $Y_xGd_{1-x}VO_4$ phases ( $x=0.25, 0.50, 0.75$ )

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### Abstract

Solid solutions of composition  $Y_xGd_{1-x}VO_4$  ( $x=0.25, 0.50, 0.75$ ) have been prepared and characterized. These phases are isomorphous with zircon,  $ZrSiO_4$ . The crystal structures were refined from X-ray powder diffraction data by the Rietveld method. The lattice parameters show a linear decrease with increasing yttrium content and the final reliability factors are in agreement with the proposed structural model. In order to study the effect produced by the substitution of  $Gd^{3+}$  by diamagnetic  $Y^{3+}$  ions, the magnetic susceptibility of  $Gd_{0.5}Y_{0.5}VO_4$  has also been measured in the temperature range 4.2–300 K.

### 1. Introduction

Rare earth ions in an oxide matrix have received considerable interest, mainly concentrated on the luminescent and magnetic properties of these ions. In fact,  $YVO_4$  activated with  $Eu^{3+}$  is a highly efficient luminophor [1].

In the case of rare earth (R) vanadates,  $RVO_4$ , all the members crystallize in the tetragonal system, space group  $I4_1/amd$ , being isomorphous with mineral zircon,  $ZrSiO_4$ . The structural parameters of  $ErVO_4$  [2],  $DyVO_4$  [3] and  $NdVO_4$  [4] have been reported previously, as well as the theoretical parameters for the other members of the series.

The rare earth vanadates (IV) have emerged as classic examples of materials which exhibit cooperative Jahn–Teller structural phase transitions [5–9]. Crystals of  $GdVO_4$  and  $YVO_4$  do not show structural phase changes [10].

In this paper we report the synthesis and structural characterization by X-ray powder diffraction of the solid solutions  $Y_xGd_{1-x}VO_4$  ( $x=0.25, 0.50, 0.75$ ). The magnetic behaviour of one member of this solid solution series has also been studied.

### 2. Experimental details

Microcrystalline materials  $Y_xGd_{1-x}VO_4$  ( $x=0.25, 0.50, 0.75$ ) were obtained by solid state reaction between the oxides  $Y_2O_3$  (Merck)  $Gd_2O_3$  (Merck)

and  $V_2O_5$ , this last oxide being isolated by the method described in ref. 11. The initial mixture in stoichiometric ratio was heated in air for 92 h at 900 K.

Powder X-ray diffraction patterns were registered at a rate of  $0.1^\circ$  ( $2\theta$ ) per minute by means of a Siemens Kristalloflex diffractometer powered by a D500 goniometer using nickel-filtered Cu  $K\alpha$  radiation.

Magnetic susceptibility measurements were made using the Faraday method in the range 4.2–300 K with a DMS5 Pendule susceptometer. The maximum applied magnetic field was 15 kG with  $H(dH/dZ) = 29 \text{ kG}^2 \text{ cm}^{-1}$ . The equipment was calibrated with  $\text{Hg}(\text{Co}(\text{SCN})_4)$  and  $\text{Gd}_2(\text{SO}_4)_3 \cdot 8\text{H}_2\text{O}$  and  $\chi$  was independent of the field in the temperature range of the measurements. Values of effective magnetic moment ( $\mu_B$ ) and Weiss constant ( $\theta$ ) were obtained by least-squares fits from the linear part of the  $\chi^{-1}$  vs.  $T$  plots.

### 3. Results and discussion

A  $2\theta$  scan of  $0.04^\circ$  was used and the Rietveld profile analysis method [12] was employed for refinement of X-ray diffraction results in the range  $10^\circ$ – $120^\circ$  ( $2\theta$ ) for 144 observed reflections. The structure of  $\text{YVO}_4$  [1] was used as trial model. The Rietveld programme minimized the function  $\chi^2 = (R_{\text{WP}}/R_{\text{EXP}})^2$ .

The reflection conditions are  $h+k+l=2n$  for  $(hkl)$ ,  $h, k=2n$  for  $(hk0)$ ,  $k+l=2n$  for  $(0kl)$ ,  $2h+l=4n$  for  $(hhl)$ ,  $l=4n$  for  $(00l)$ ,  $h=2n$  for  $(h00)$  and  $h=2n$  for  $(h\bar{h}0)$ , compatible with the space group  $I4_1/amd$  ( $Z=4$ ). The atoms occupy the following positions: yttrium and gadolinium in  $4a$  ( $0, \frac{3}{4}, \frac{1}{8}$ ), vanadium in  $4b$  ( $0, \frac{1}{4}, \frac{3}{8}$ ) and oxygen in  $16h$  ( $0, y, z$ ).

The crystallographic parameters are presented in Table 1. As can be seen, the unit cell parameters decrease linearly with increasing the yttrium content.

The structure of these solids solutions is isomorphous with zircon [14] and can be described as chains of alternating  $\text{RO}_8$  distorted polyhedra (with  $R \equiv \text{Y}$  or  $\text{Gd}$  randomly distributed) sharing edges and corners with  $\text{VO}_4$

TABLE 1  
Crystallographic parameters of  $\text{Y}_x\text{Gd}_{1-x}\text{VO}_4$

Composition	Unit cell parameters (Å)		Oxygen position		Reliability factors	
	<i>a</i>	<i>c</i>	<i>y</i>	<i>z</i>	$R_{\text{WP}}$	$R_{\text{P}}$
$\text{GdVO}_4$ [13]	7.213	6.348				
$\text{Y}_{0.25}\text{Gd}_{0.75}\text{VO}_4$	7.164	6.315	0.435	0.212	13.52	8.22
$\text{Y}_{0.5}\text{Gd}_{0.5}\text{VO}_4$	7.143	6.310	0.440	0.201	12.50	7.19
$\text{Y}_{0.75}\text{Gd}_{0.25}\text{VO}_4$	7.136	6.300	0.436	0.202	14.68	5.31
$\text{YVO}_4$ [1]	7.123	6.292				

tetrahedra. Figure 1 shows a perspective of the structure in which chains of  $\text{VO}_4(1)\text{-RO}_8\text{-VO}_4(1)$  are defined along the  $c$  axis direction. It can be seen that there is octa-coordination of R ions (black oxygen atoms), with  $\text{VO}_4$  tetrahedra offering a common edge (labelled 1) or a vertex (labelled 2 and 3). This fact is reflected in two sets of R-O bond distances, whilst the V-O ones remain unchanged with the composition (Table 2).

The magnetic susceptibility of  $\text{Y}_{0.5}\text{Gd}_{0.5}\text{VO}_4$  has also been measured in the temperature range 4.2–300 K as an example. These results show a quasi-Curie-like behaviour over the entire interval of measurements, as can be observed from Figure 2. This behaviour is characteristic of the ground state of  $\text{Gd}^{3+}$ ,  $^8\text{S}_{7/2}$ , for which the crystal field splitting should be very small, *i.e.* less than  $1\text{ cm}^{-1}$  [15], and can be neglected, and the magnetic moment takes the experimental value  $8.01\ \mu_{\text{B}}$ , which agrees with the free-ion value of  $7.94\ \mu_{\text{B}}$  [16].

The high symmetry of the ground state of  $\text{Gd}^{3+}$  yields a  $\theta$  value for  $\text{Y}_{0.5}\text{Gd}_{0.5}\text{VO}_4$  close to zero ( $\theta = -0.42$ ). It is obvious that this negative Weiss constant value is due to the crystal field effects and non-cooperative interactions, *i.e.* antiferromagnetic coupling, present in this compound down to 4.2 K.

By substitution of  $\text{Gd}^{3+}$  by  $\text{Y}^{3+}$  ions the number of paramagnetic ions is reduced but the magnetic behaviour is scarcely modified with respect to  $\text{GdVO}_4$  [11]. Analogous results are obtained for the remaining stoichiometries

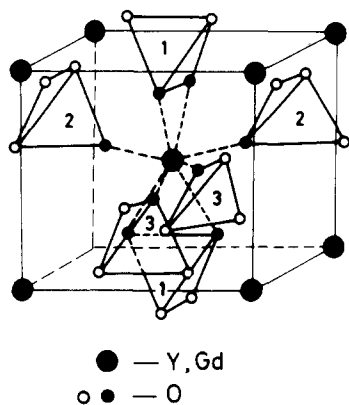


Fig. 1.  $\text{VO}_4$  tetrahedra around yttrium or gadolinium atoms.

TABLE 2

Selected interatomic distances (angstroms) in  $\text{Y}_x\text{Gd}_{1-x}\text{VO}_4$  ( $\text{R} \equiv \text{Y}$  or  $\text{Gd}$ )

	$\text{Y}_{0.25}\text{Gd}_{0.75}\text{VO}_4$	$\text{Y}_{0.5}\text{Gd}_{0.5}\text{VO}_4$	$\text{Y}_{0.75}\text{Gd}_{0.25}\text{VO}_4$
R-O ( $\times 4$ )	2.29(3)	2.28(3)	2.31(9)
R-O ( $\times 4$ )	2.45(4)	2.50(8)	2.51(9)
V-O ( $\times 4$ )	1.71(5)	1.71(8)	1.71(9)

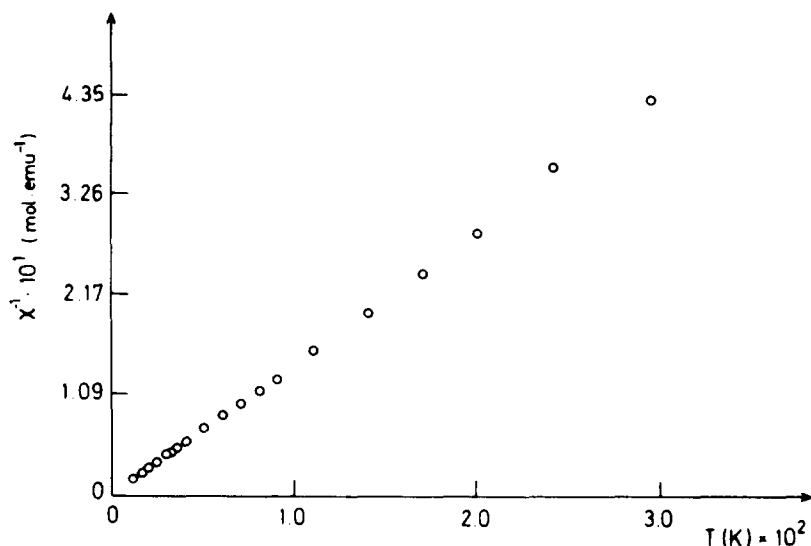


Fig. 2. Temperature dependence of reciprocal molar susceptibility in  $Y_{0.5}Gd_{0.5}VO_4$ .

in the solid solutions. From the crystal structure of these oxides one can conclude that in the whole range of composition paramagnetic ions are isolated, taking into account their location on  $4a$  crystallographic sites. Therefore it is feasible to modulate the  $Gd^{3+}$  concentration in the material for different purposes.

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